

Algunas cuestiones sobre la resolución de orden reducido de ecuaciones en derivadas parciales

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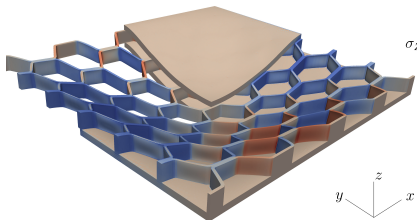
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Modelización de orden reducido

- Hoy en día muchos problemas en la ciencia y la ingeniería siguen siendo intratables, a pesar de los impresionantes avances recientes en modelado, análisis numérico, técnicas de discretización y computación.
 - Por ejemplo en química cuántica o flujo de gases enrarecidos, los modelos matemáticos están planteados en espacios de dimensión alta (D). Utilizando una malla estándar con $M \simeq 10^3$ nodos, si $D \simeq 30$ (un modelo muy simple), el número de grados de libertad es cercano a 10^{90} !!.
- La optimización de sistemas, los problemas inversos y la cuantificación de la incertidumbre son muy costosas computacionalmente, al requerir el cálculo reiterado del estado del sistema, en función de los parámetros de diseño. Esto necesita horas, días y semanas de computación. Soluciones en tiempo real o abordable para el diseño están frecuentemente fuera de alcance.

Modelización de orden reducido

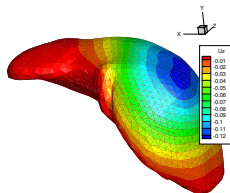
- La modelización de orden reducido (ROM) se basa en la construcción de variedades de dimensión baja que aproximen bien la familia de soluciones paramétricas buscadas.
- Permite conseguir reducciones dramáticas del tiempo de cálculo.
- Las matemáticas son cruciales en la elaboración de ROMs eficaces.



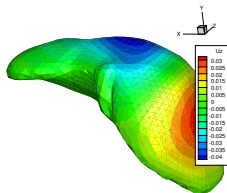
- Modelos reducidos de turbulencia (Modelo de Smagorinsky). Técnica de bases reducidas.
 - Aplicación al análisis del comportamiento energético de patios.
- Modelización de orden reducido de problemas elípticos paramétricos. Técnica PGD (Proper Generalized Decomposition).

Un ejemplo: Cirugía virtual del hígado

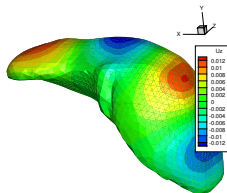
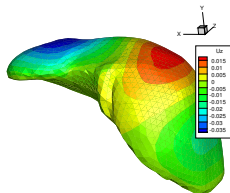
- Modos propios de vibración elástica de un hígado paramétrico.



(a)



(b)



Application to Smagorinsky turbulence model

- Ph. D. of Enrique Delgado, SINUM (2018).

Start from Navier-Stokes equations: Let $\Omega \subset \mathbb{R}^d$ bounded, with

$\partial\Omega = \Gamma = \Gamma_D \cup \Gamma_N$, where $\Gamma_D = \Gamma_{D_g} \cup \Gamma_{D_0}$

$$\left\{ \begin{array}{ll} \mathbf{u} \cdot \nabla \mathbf{u} - \frac{1}{\mu} \Delta \mathbf{u} + \nabla p = \mathbf{f} & \text{in } \Omega \\ \nabla \cdot \mathbf{u} = 0 & \text{in } \Omega \\ \mathbf{u} = \mathbf{g}_D & \text{on } \Gamma_{D_g} \\ \mathbf{u} = 0 & \text{on } \Gamma_{D_0} \\ -p\mathbf{n} + \left(\frac{1}{\mu} \right) \frac{\partial \mathbf{u}}{\partial \mathbf{n}} = 0 & \text{on } \Gamma_N \end{array} \right. \quad (1)$$

Smagorinsky model

$$\left\{ \begin{array}{ll} \mathbf{w} \cdot \nabla \mathbf{w} - \frac{1}{\mu} \Delta \mathbf{w} + \nabla p - \nabla \cdot (\nu_T(\mathbf{w}) \nabla \mathbf{w}) = \mathbf{f} & \text{in } \Omega \\ \nabla \cdot \mathbf{w} = 0 & \text{in } \Omega \\ \mathbf{w} = \mathbf{g}_D & \text{on } \Gamma_{D_{in}} \\ \mathbf{w} = 0 & \text{on } \Gamma_{D_w} \\ -p\mathbf{n} + \left(\frac{1}{\mu} + \nu_T(\mathbf{w}) \right) \frac{\partial \mathbf{w}}{\partial \mathbf{n}} = 0 & \text{on } \Gamma_{out} \end{array} \right. \quad (2)$$

where $\nu_T(\mathbf{w})(x) = (C_S h_K)^2 |\nabla \mathbf{w}|_K(x)|$.

Model intrinsically discrete, linked to a discretization grid.

LPS-Smagorinsky model

Finite element spaces

$$Y_h = \overline{Y}_h \oplus Y'_h \quad M_h = \overline{M}_h \oplus M'_h, \quad \sigma_h : Y_h \mapsto \overline{Y}_h : \text{Averaging operator}$$

$$\text{thus, } \mathbf{u}_h = \overline{\mathbf{u}}_h + \mathbf{u}'_h, \quad \mathbf{u}'_h = (Id - \sigma_h)\mathbf{u}_h = \sigma_h^* \mathbf{u}_h, \quad p_h = \overline{p}_h + p'_h$$

LES Closure model: Smagorinsky eddy diffusion,

$$\nu_t(\mathbf{u}_h) = (C_S h_K)^2 |\nabla(\mathbf{u}_h)| \text{ on element } K \in \mathcal{T}_h,$$

$$a_s(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h) = \int_{\Omega} \nu_t(\mathbf{u}'_h) \nabla \mathbf{u}'_h : \nabla \mathbf{v}'_h d\Omega$$

Pressure stabilization:

$$s_{pres}(p, q) = \int_{\Omega} \tau_{K,p}(\mu) \sigma_h^*(\nabla p_h) \sigma_h^*(\nabla q_h) d\Omega,$$

$$\text{with stabilization coefficients } \tau_{K,p}(\mu) = \left[c_1 \frac{1/\mu + \overline{\nu_T}}{h_K^2} + c_2 \frac{U_K}{h_K} \right]^{-1}.$$

Full Order problem

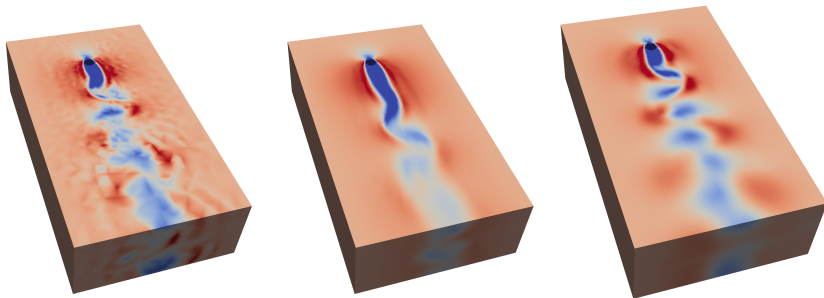
LPS-Smagorinsky model

$$\left\{ \begin{array}{ll} \text{Find } (\mathbf{u}_h, p_h) \in X_h \text{ such that} \\ a(\mathbf{u}_h, \mathbf{v}_h; \mu) + b(\mathbf{v}_h, p_h; \mu) + c(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h; \mu) \\ + a_s(\mathbf{u}_h; \mathbf{u}_h, \mathbf{v}_h; \mu) = \langle f, \mathbf{v}_h \rangle & \forall \mathbf{v}_h \in Y_h \\ b(\mathbf{u}_h, q_h; \mu) + s_{pres}(p_h, q_h; \mu) = 0 & \forall q_h \in M_h, \end{array} \right. \quad (3)$$

where the operator terms are

a : Diffusion, b : Pressure gradient-Divergence, c : Convection,
 a_s Eddy diffusion, s_{pres} : Pressure discretization stabilization.

Effect of LPS stabilization for convection: Flow past a cylinder



- Quadratic elements with 4×26.512 degrees of freedom

Magnitude of the velocity for $Re = 200$.

Convection stabilization by

- Plain Galerkin method (left): No stabilization
- Pure penalty method (center): $\sigma_h^* = Id$.
- Projection-stabilized one (right): $\sigma_h^* = Id - \sigma_h$.

Construction of Reduced Basis Space: Reduced basis problem

Given the space $X_N = Y_N \times M_N, 1 \leq N \leq N_{\max}$

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$$A(U_N(\mu), V_N; \mu) = F(V_N) \quad \forall V_N \in X_N$$

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- Choose μ_{N+1} as

$$\mu_{N+1} = \arg \max_{\mu \in \overline{\mathcal{D}}} \|U_h(\mu) - U_N(\mu)\|_X$$

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- Construct the new reduced spaces

$$Y_{N+1} = \text{span}\{\zeta_i^y := \mathbf{u}(\mu^i), i = 1, \dots, N+1\}$$

$$M_{N+1} = \text{span}\{\zeta_i^p := p(\mu^i), i = 1, \dots, N+1\}$$

Reduced Basis problem

Given the space $X_N = Y_N \times M_N$, $1 \leq N \leq N_{\max}$,

$$\left\{ \begin{array}{ll} \text{Find } (\mathbf{u}_N, p_N) \in X_N \text{ such that} & \\ a(\mathbf{u}_N, \mathbf{v}_N; \mu) + b(\mathbf{v}_N, p_N; \mu) + c(\mathbf{u}_N; \mathbf{u}_N, \mathbf{v}_N; \mu) & \\ + a_s(\mathbf{u}_N; \mathbf{u}_N, \mathbf{v}_N; \mu) = \langle f, \mathbf{v}_N \rangle & \forall \mathbf{v}_N \in Y_N \\ b(\mathbf{u}_N, q_N; \mu) + s_{pres}(p_N, q_N; \mu) = 0 & \forall q_N \in M_N, \end{array} \right.$$

where we recall the operator terms:

a : Diffusion, b : Pressure gradient-Divergence, c : Convection,
 a_s Eddy diffusion, s_{pres} : Pressure discretization stabilization.

- Typically, $\dim(X_M)$ is much smaller than $\dim(X_h)$.

Empirical interpolation of turbulent viscosity and stabilized coefficients

- The eddy diffusion term is approximated as

$$g(x, \mu) = \nu_t(\mathbf{u}_N)(x, \mu) \simeq \sum_{j=1}^M \alpha_j(\mu) \mathbf{q}_j(x),$$

where $\{\mathbf{q}_1, \dots, \mathbf{q}_M\}$ are a linear combination of particular snapshots $g(x, \mu^1), \dots, g(x, \mu^M)$ determined, together with the interpolation points x_i , hierarchically to improve the approximation properties of interpolation operator by incorporating iteratively the worse case.

The coefficients $\alpha_j(\mu)$ are determined to interpolate g at the x_i ,

$$\sum_{j=1}^M \alpha_j(\mu) \mathbf{q}_j(x_i) = g(x_i, \mu) \quad i = 1, \dots, M$$

- A similar approximation is built for the stabilized coefficients $\tau_{K,p}(\mu)$.

Role of mathematics: A-posteriori error estimator

- It holds

$$|\partial_1 A(U^1, V; \mu)(Z) - \partial_1 A(U^2, V; \mu)(Z)| \leq \rho_T \|U^1 - U^2\|_X \|Z\|_X \|V\|_X.$$

$$\text{Let } \beta_N(\mu) = \inf_{Z_h \in X_h} \sup_{V_h \in X_h} \frac{\partial_1 A(U_N(\mu), V_h; \mu)(Z_h)}{\|Z_h\|_X \|V_h\|_X}, \quad \tau_N(\mu) = \frac{4\epsilon_N(\mu)\rho_T}{\beta_N^2},$$

$$\text{with } \epsilon_N(\mu) = \|\mathcal{R}(\cdot; \mu)\|_{X'}, \quad \mathcal{R}(V_h; \mu) = F(V_h; \mu) - A(U_N(\mu), V_h; \mu)$$

Theorem

If $\beta_N > 0$ and $\tau_N(\mu) \leq 1$, then there exists a unique solution $U_h(\mu)$ to (FE) such that

$$\|U_h(\mu) - U_N(\mu)\|_X \leq \Delta_N(\mu),$$

$$\text{where } \Delta_N(\mu) = \frac{\beta_N}{2\rho_T} \left[1 - \sqrt{1 - \tau_N(\mu)} \right].$$

Reynolds range: $\mu \in [1000, 5100]$

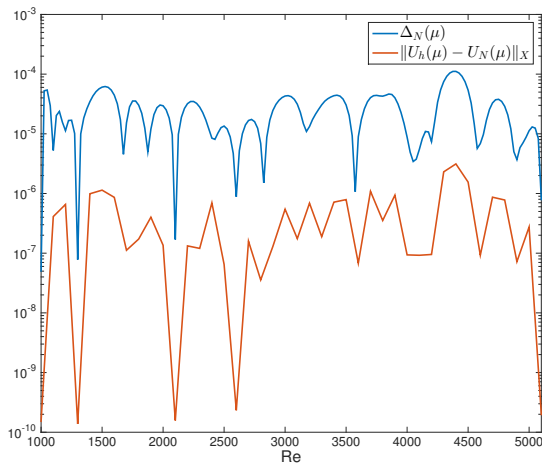


Figure: *A posteriori* error bound at $N=16$

FE and RB velocity solution

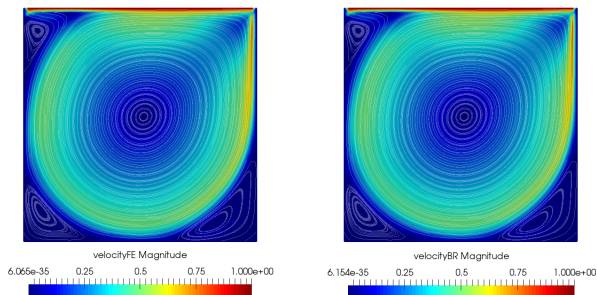


Figure: FE (left) and RB (right) velocity solution for $\mu = 4521$

Results: Error and speed-up analysis.

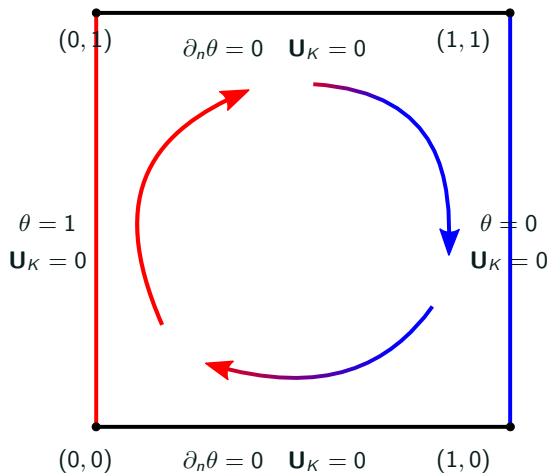
FE dof: 30603

EIM dof: $25 (\nu_T) + 20 (\tau_{K,p})$, RB dof: 32

Data	$\mu = 1610$	$\mu = 2751$	$\mu = 3886$	$\mu = 4521$
T_{FE}	4083.19s	6918.53s	9278.51s	10201.7s
T_{online}	0.71s	0.69s	0.69s	0.7s
speedup	5750	10026	13280	14459
$\ \mathbf{u}_h - \mathbf{u}_N\ _T$	$2.4 \cdot 10^{-5}$	$4.129 \cdot 10^{-6}$	$3.14 \cdot 10^{-5}$	$3.23 \cdot 10^{-5}$
$\ p_h - p_N\ _0$	$2.17 \cdot 10^{-7}$	$1.99 \cdot 10^{-8}$	$5.38 \cdot 10^{-8}$	$6.36 \cdot 10^{-8}$

Table: Data summary

Boussinesq-Smagorinsky model. Problem description



Boussinesq-Smagorinsky model

$$\left\{ \begin{array}{ll} \mathbf{U}_K \cdot \nabla \mathbf{U}_K - Pr \Delta \mathbf{U}_K + \nabla p - \nabla \cdot (\nu_T(\mathbf{U}_K) \nabla \mathbf{U}_K) = \mathbf{f} + Pr Ra \theta \mathbf{e}_2 & \text{in } \Omega \\ \nabla \cdot \mathbf{U}_K = 0 & \text{in } \Omega \\ \mathbf{U}_K \cdot \nabla \theta - \Delta \theta - \frac{1}{Pr} \nabla \cdot (\nu_T(\mathbf{U}_K) \nabla \theta) = Q & \text{in } \Omega \\ \mathbf{U}_K = 0 & \text{on } \Gamma \\ \theta = \theta_D & \text{on } \Gamma_D \\ \theta = 0 & \text{on } \Gamma_0 \\ \partial_n \theta = 0 & \text{on } \Gamma_N \end{array} \right.$$

where, $\nu_T(\mathbf{U}_K)(x) = (C_S h_K)^2 |\nabla(\Pi_h \mathbf{U}_K)|_K(x)|$

FE velocity-temperature solution

- Rayleigh range: $Ra \in [10^3, 10^5]$.
- Taylor-Hood Finite Element, $(\mathbb{P}2 - \mathbb{P}2 - \mathbb{P}1)$, Regular mesh (2601 nodes and 5000 triangles).

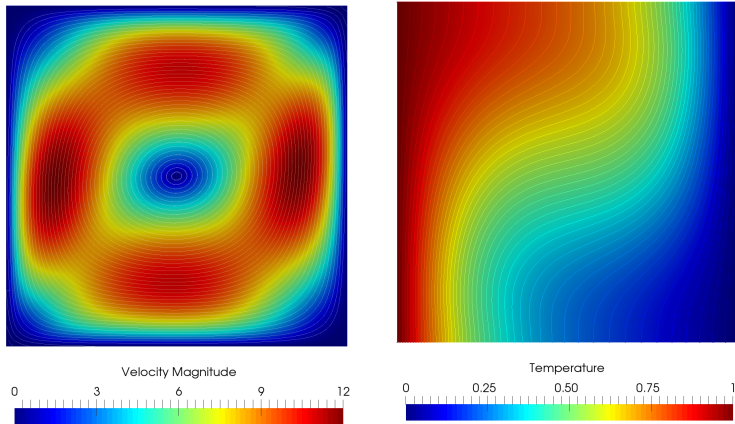


Figure: FE velocity and temperature solution for $Ra = 4363$

FE velocity-temperature solution

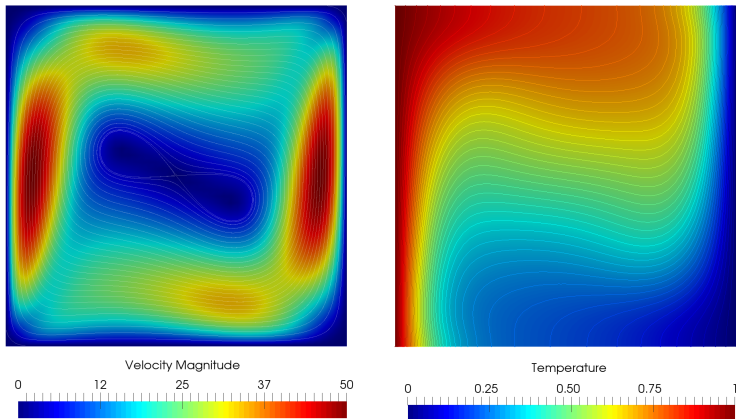


Figure: FE velocity and temperature solution for $Ra = 53778$

FE velocity-temperature solution. Large Rayleigh range: $Ra \in [10^5, 10^6]$

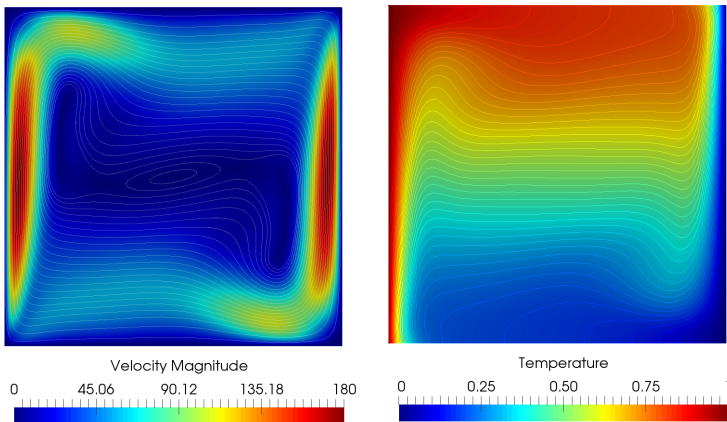


Figure: FE velocity and temperature solution for $Ra = 667746$

Results: Error and speed-up analysis.

- Moderate Rayleigh number range, $Ra \in [10^3, 10^5]$:

FE dof: 33204, EIM dof: 42, RB dof: 88.

Data	$Ra = 4060$	$Ra = 17808$	$Ra = 53778$	$Ra = 93692$
T_{FE}	633.65s	585.83s	553.25s	677.86s
T_{online}	0.55s	0.5s	0.46s	0.49s
speedup	1133	1151	1189	1367

- Large Rayleigh number range, $Ra \in [10^5, 10^6]$:

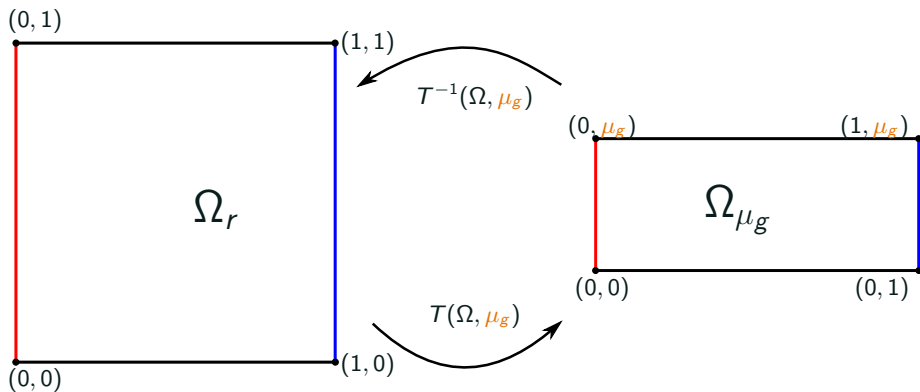
FE dof: 64684, EIM dof: 150, RB dof: 256

Data	$\mu = 16941$	$\mu = 355402$	$\mu = 667746$	$\mu = 921441$
T_{FE}	3563.11s	3675.01s	4354.26s	4928.37s
T_{online}	9.28s	11.34s	15.22s	16.8s
speedup	383	324	285	293

- Errors in H^1 norm below 10^{-5} in all cases.

Geometrical parameters

- Change of variable from reference domain.



$$DT(\Omega; \mu_g) = \begin{pmatrix} 1 & 0 \\ 0 & \mu_g \end{pmatrix}.$$

Variational formulation

Boussinesq-Smagorinsky F.E. with geometrical parametrization

$$\left\{ \begin{array}{l} \text{Given } \boldsymbol{\mu} = (\mu_g, Ra) \in \mathcal{D} = \mathcal{D}_g \times \mathcal{D}_{Ra}, \\ \text{find } U_h(\boldsymbol{\mu}) = (\mathbf{u}_h, p_h, \theta_h) \in X_h \text{ s.t.} \\ A(U_h(\boldsymbol{\mu}), V_h; \mu_g) = F(V_h; Ra) \quad \forall V_h \in X_h \end{array} \right. \quad (4)$$

- Same finite element spaces for all parameters.

FE velocity solutions. Parameter range:
 $Ra = 10^5, \mu_g \in [0.5, 2]$.

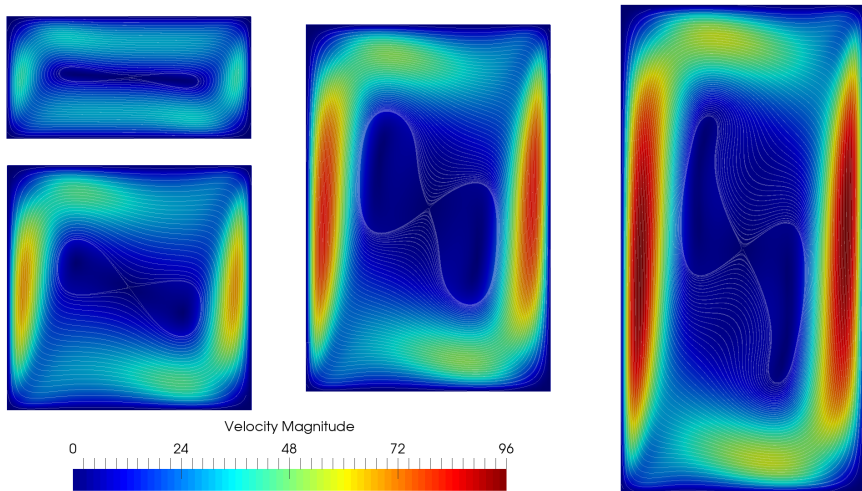
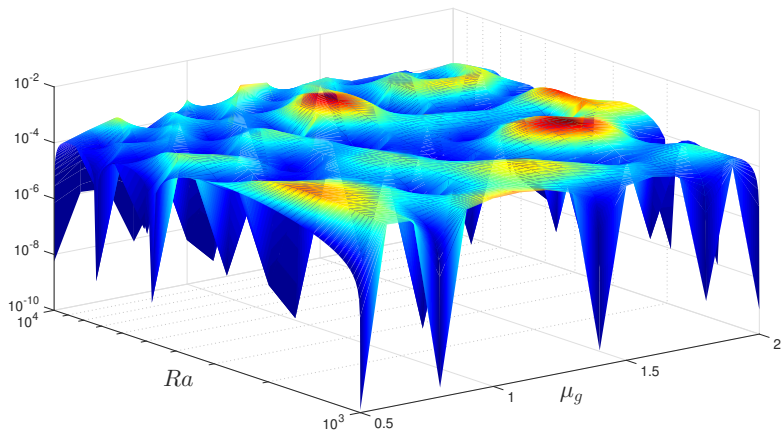


Figure: FE velocities for different values of μ_g . ($Ra = 10^5$)

A posteriori error bound at $N = N_{\max}$



Results: Error and speed-up analysis.

- Geometric parameter, $\mu \in [0.5, 2]$, $Ra = 10^5$

FE dof: 33204, EIM dof: 73, RB dof: 128 $\mu \in [0.5, 2]$

Data	$\mu_g = 0.64$	$\mu_g = 1.08$	$\mu_g = 1.44$	$\mu_g = 1.87$
T_{FE}	808.91s	810.16s	866.1s	851.82s
T_{online}	2.68s	2.55s	2.61s	2.52s
speedup	301	317	331	337

- Geometric and physical parameter, $\mu \in [0.5, 2]$, $Ra \in [10^3, 10^4]$ FE dof: 33204, EIM dof: 138, RB dof: 216

Data	$Ra = 2143$ $\mu_g = 1.95$	$Ra = 3506$ $\mu_g = 0.71$	$Ra = 5922$ $\mu_g = 1.13$	$Ra = 9618$ $\mu_g = 1.63$
T_{FE}	600.96s	914.18s	684.95s	630.94s
T_{online}	11.08s	15.73s	14.52s	11.46s
speedup	54	58	47	55

Present and future work

- **Turbulence modeling:**

- Extension to transient and 3D flows.
- Application to more complex turbulence models (VMS).
- Applications to energy-efficient design of buildings.

- **Numerical Analysis:**

- Take advantage of VMS formulation: Use the modeled u' as error indicator.
- Further analysis of inf-sup conditions for RBM.
- Error estimates for EIM of eddy viscosity.

Parametric elliptic problems

Let

- A separable Hilbert space $(H, (\cdot, \cdot))$.
- A measure space $(\Gamma, \mathcal{B}, \mu)$, with standard notations, so that μ is σ -finite.
- A form $a \in L^\infty(\Gamma, B_s(H); d\mu)$ uniformly elliptic and coercive on H w. r. t. γ .
- A data function $f \in L^2(\Gamma, H'; d\mu)$

We are interested in solving the variational problem:

Find $u(\gamma) \in H$ such that

$$a(u(\gamma), v; \gamma) = \langle f(\gamma), v \rangle_{H'-H}, \quad \forall v \in H, \text{ } d\mu\text{-a.e. } \gamma \in \Gamma, \quad (5)$$

This is a common situation in many engineering problems when the properties of the media are parameter-depending.

Proper Orthogonal Decomposition

Objective: Approximate

$$u(\gamma) \simeq \sum_{k \geq 1} \Phi_k(\gamma) w_k, \quad w_k \in H.$$

The POD searches for $w_1, w_2, \dots$ such that for each $n = 1, 2, \dots$ the space $S_n = \text{Span}\{w_1, \dots, w_n\}$ minimizes

$$\int_{\Gamma} \|u(\gamma) - u_Z(\gamma)\|_H^2 d\mu(\gamma)$$

among all sub-spaces Z of H of dimension n , where $u_Z(\gamma) \in Z$ solves

$$a(u_Z(\gamma), z; \gamma) = \langle f(\gamma), z \rangle, \quad \forall z \in Z, \quad d\mu\text{-a.e. } \gamma \in \Gamma, \quad (6)$$

Proper Orthogonal Decomposition

The $w_1, w_2, \dots$ turn out to be the eigenfunctions of the POD operator $\mathcal{R} : H \mapsto H$, given by

$$\mathcal{R}(v) = \int_{\Gamma} (v, u(\gamma))_H u(\gamma) d\mu(\gamma), \text{ for } v \in H.$$

This operator is self-adjoint, positive and compact, what gives the existence (and uniqueness) of the optimal sub-spaces.

- The POD is a generalization of the Singular Values Decomposition of matrices. Also called Principal Component Analysis, Karhunen-Loève decomposition,
- Needs to know the inner products $(v, u(\gamma))_H$ to compute $\mathcal{R}$. In practice a quadrature formula is applied and only a certain number of $u(\gamma_i)$ (“snapshots”) need to be computed.
- Then the targeted PDE needs to be a-priori solved to compute the POD expansion of its parametric solution.

Proper Generalized Decomposition

The PGD searches also for a similar tensorized decomposition

$$u(\gamma) \simeq \sum_{k \geq 1} \Phi_k(\gamma) w_k, \quad u_k \in H.$$

but the w_k are computed online, iteratively.

- This is done by partial Galerkin problems.
- For instance, the first summand $u_1(\gamma) = \Phi_1(\gamma) w_1$, with $\Phi_1(\gamma) \in L^2(\Gamma, d\mu)$ and $w_1 \in H$ is a solution of

$$\begin{aligned} a(\Phi_1(\gamma) w_1, \Phi_1(\gamma) v) &= \langle f(\gamma), \Phi_1(\gamma) v \rangle, \quad \forall v \in H, d\mu\text{-a.e. } \gamma \in \Gamma; \\ \int_{\Gamma} a(\Phi_1(\gamma) w_1, w_1) s(\gamma) d\mu(\gamma) &= \int_{\Gamma} \langle f(\gamma), w_1 \rangle s(\gamma) d\mu(\gamma), \quad \forall s \in L^2(\Gamma, d\mu). \end{aligned}$$

- These problems are solved by a power-iteration algorithm

Proper Generalized Decomposition

The following term

$$u_k(\gamma) = \sum_{k=1}^n \Phi_k(\gamma) w_k, = u_{k-1} + \Phi_{k-1}(\gamma) w_{k-1} \quad w_{k-1} \in H.$$

is computed by a deflation algorithm, i. e., the same but replacing f by the current residual, $r_{k-1}(\gamma) = f(\gamma) - A(\gamma)u_{k-1}$.

- The PGD has been characterized by a non-optimal descent method for elliptic problems (Falcó and Nouy, 2016).

Optimal sub-spaces of finite dimension

Problem targeted: Find the best subspace W of H of dimension $\leq k$ that minimizes the mean quadratic error between $u(\gamma)$ and $u_W(\gamma)$ with respect to the norm generated by the form $a(\cdot, \cdot; \gamma)$.

That is, W solves

$$(P) \quad \min_{Z \in \mathcal{S}_k} \int_{\Gamma} a(u(\gamma) - u_Z(\gamma), u(\gamma) - u_Z(\gamma); \gamma) d\mu(\gamma), \quad (7)$$

where $\mathcal{S}_k$ is the family of subspaces of H of dimension $\leq k$ and $u_Z(\gamma) \in Z$ is the solution of the Galerkin approximation of problem (6) on Z ,

$$a(u_Z(\gamma), z; \gamma) = \langle f(\gamma), z \rangle, \quad \forall z \in Z, \quad d\mu\text{-a.e. } \gamma \in \Gamma,$$

A look at the 1D case

When $k = 1$, problem (P) can be written as

$$\min_{v \in H, \varphi \in L^2(\Gamma; d\mu)} \int_{\Gamma} a(u(\gamma) - \varphi(\gamma)v, u(\gamma) - \varphi(\gamma)v; \gamma) d\mu(\gamma).$$

So, taking the derivative of the functional

$$(v, \varphi) \in H \times L^2(\Gamma; d\mu) \mapsto J(v, \varphi) = \int_{\Gamma} a(u(\gamma) - \varphi(\gamma)v, u(\gamma) - \varphi(\gamma)v; \gamma) d\mu(\gamma),$$

we deduce that w is a solution of the non-linear variational problem

$$\int_{\Gamma} \frac{a(u(\gamma), w; \gamma)}{a(w, w, \gamma)} a(u(\gamma), v; \gamma) d\mu(\gamma) = \int_{\Gamma} \frac{a(u(\gamma), w; \gamma)^2}{a(w, w, \gamma)^2} a(w, v; \gamma) d\mu(\gamma), \quad (8)$$

$$\forall v \in H.$$

A look at the 1D case

- If a does not depend on γ , statement (8) is equivalent to

$$\mathcal{R}w = \int_{\Gamma} a(u(\gamma), w) u(\gamma) d\mu(\gamma) = \lambda w,$$

where

$$\lambda = \frac{\int_{\Gamma} a(u(\gamma), w)^2 d\mu(\gamma)}{a(w, w)}.$$

i.e. w is an eigenvector of the POD operator $\mathcal{R}$ when the inner product in H is the form $a(\cdot, \cdot)$.

- However, when a depends on γ problem (8) does not correspond to a proper eigenvalue equation: It is a non-linear eigenfunction problem, where no eigenvalues appear.

Then, we are considering a genuine extension of the POD.

Optimal sub-spaces

Theorem

There exists at least a sub-space $Z \in S_k$ that solves problem (P)

- Proof by direct method of the Calculus of Variations + compactness argument.
- Special proof for the 1D case. This problem is equivalent to

$$(P') \quad \max_{\substack{\Psi \in H \\ \|\Psi\|=1}} \int_{\Gamma} \frac{\langle f(\gamma), \Psi \rangle^2}{a(\Psi, \Psi; \gamma)} d\mu(\gamma).$$

Theorem

Problem (P') admits at least one solution.

- Proof by combination of direct method of the Calculus of Variations, compactness in Hilbert spaces, uniform boundedness and ellipticity of forms $a(\cdot, \cdot; \gamma)$ and Fatou's Lemma.

Characterization of PGD algorithm

- The problem to compute the PGD modes,

$$\begin{aligned} a(\Phi_1(\gamma) w_1, \Phi_1(\gamma) v) &= \langle f(\gamma), \Phi_1(\gamma) v \rangle, \quad \forall v \in H, d\mu\text{-a.e. } \gamma \in \Gamma; \\ \int_{\Gamma} a(\Phi_1(\gamma) w_1, w_1) s(\gamma) d\mu(\gamma) &= \int_{\Gamma} \langle f(\gamma), w_1 \rangle s(\gamma) d\mu(\gamma), \quad \forall s \in L^2(\Gamma, d\mu). \end{aligned}$$

turns out to be the first-order optimality conditions of the critical points of the functional $J(v, \varphi)$.

Tensor approximation

Similarly to the PGD, we expand $u(\gamma)$ by the tensor approximation

$$u(\gamma) = \sum_{k \geq 1} \Phi_k(\gamma) w_k, \quad w_k \in H.$$

where the w_k are obtained by a deflation algorithm similar to the one followed by PGD:

- Initialization:

$$w_1 = \operatorname{argmin}_{\psi \in H} \int_{\Gamma} a(u(\gamma) - u_{\psi}(\gamma), u(\gamma) - u_{\psi}(\gamma); \gamma) d\mu(\gamma),$$

where u_{ψ} is the Galerkin solution of the targeted elliptic problem on $\operatorname{span}\{\psi\}$.

- Iteration: Known $u_{k-1}(\gamma) = \sum_{i=1}^{k-1} \Phi_i(\gamma) w_i$, let $e_{k-1} = u - u_{k-1}$.

$$w_k = \operatorname{argmin}_{\psi \in H} \int_{\Gamma} a(e_{k-1}(\gamma) - u_{\psi}(\gamma), e_{k-1}(\gamma) - u_{\psi}(\gamma); \gamma) d\mu(\gamma),$$

Tensor approximation

- It holds that $w_k = (e_{k-1})_W$, with W a solution of

$$\max_{\Psi \in H} \left\{ \int_{\Gamma} \langle f(\gamma), (e_{k-1})_{\Psi}(\gamma) \rangle d\mu(\gamma) - \bar{a}(u_{k-1}, (e_{k-1})_{\Psi}) \right\},$$

and $(e_{k-1})_{\Psi} \in L^2(\Gamma, Z; d\mu)$ the solution of

$$\bar{a}((e_{k-1})_{\Psi}, z) = \int_{\Gamma} \langle f(\gamma), z(\gamma) \rangle d\mu(\gamma) - \bar{a}(u_{k-1}, z), \quad \forall z \in L^2(\Gamma, Z; d\mu).$$

- This allows us to carry out the different iterations without needing to know the function u .

Theorem

The approximations u_n provided by the deflation algorithm strongly converge to the solution u of problem (P).

- Proof by orthogonality properties of residuals, consequence of the symmetry of forms a .

Test case

- Consider the elliptic problem with variable diffusion

$$\begin{cases} -\nabla \cdot (\mu(\gamma) \nabla u) &= f & \text{in } \Omega, \\ u &= 0 & \text{on } \partial\Omega; \end{cases}$$

where $\Omega = (0, 1)^2$ and

$$\mu(\gamma)(x, y) = \begin{cases} \gamma + \sigma & \text{if } 0 \leq x \leq 1/4, \\ 1 + \sigma & \text{if } 1/4 \leq x \leq 1. \end{cases} \quad \text{for all } (x, y) \in \overline{\Omega}.$$

σ is a real number to be selected to vary the minimum α of $\mu(\gamma)$.

Computation of PGD modes

- The PGD modes are solved by the Power Iteration Algorithm with normalization to solve for the critical points of the functional to be minimized $J(v, \varphi)$:

(a) $\tilde{w}^{n+1} \in H$ satisfying, $\forall v \in H, d\mu$ -a.e. $\gamma \in \Gamma$,

$$\int_{\Gamma} a(\Phi^n(\gamma) \tilde{w}^{n+1}, \Phi^n(\gamma) v) d\mu(\gamma) = \int_{\Gamma} \langle f(\gamma), \Phi^n(\gamma) v \rangle, d\mu(\gamma);$$

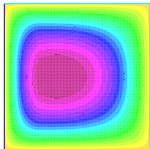
(b) $w^{n+1} = \frac{\tilde{w}^{n+1}}{\|\tilde{w}^{n+1}\|_H};$

(c) $\phi^{n+1} \in L^2(\Gamma, d\mu)$ satisfying, $\forall s \in L^2(\Gamma, d\mu),$

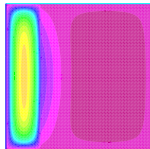
$$\int_{\Gamma} a(\Phi^{n+1}(\gamma) w^{n+1}, w^{n+1}) s(\gamma) d\mu(\gamma) = \int_{\Gamma} \langle f(\gamma), w^{n+1} \rangle s(\gamma) d\mu(\gamma), .$$

- This algorithm converges with linear rate if f/α is small enough.

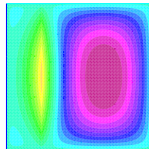
Computed modes



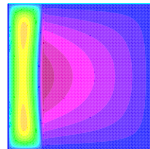
(a) Mode 1



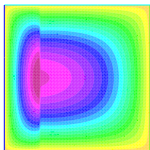
(b) Mode 2



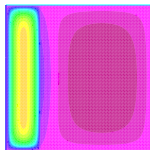
(c) Mode 3



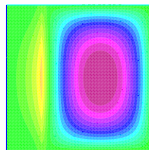
(d) Mode 4



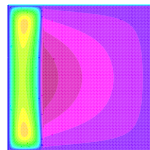
(e) Mode 5



(f) Mode 6



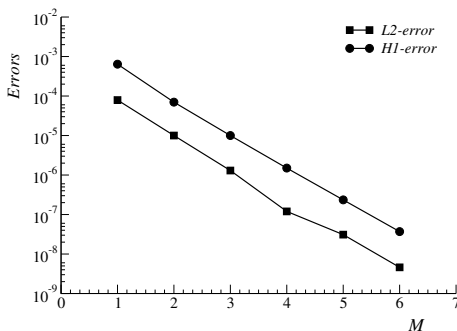
(g) Mode 7



(h) Mode 8

Convergence history

The next figure displays $\|u(\gamma) - \sum_{i=1}^M \Phi_i(\gamma) w_i\|_X$
with $X = L^2(\Gamma, L^2(\Omega); d\mu)$ and $X = L^2(\Gamma, H^1(\Omega); d\mu)$.



(i) CV

- The expansion converges with *spectral* rate, like ρ^{-M} for $\rho > 1$.

Concluding remarks

- We have constructed an on-line tensorized approximation of parameterized elliptic equations (similar to PGD) with optimal approximation of each summand and orthogonality between summands (similar to POD).
- Optimization algorithms to compute the modes.
- Extensions to evolution and non-linear problems.
- Extension to more complex basic approximation structures (tensor trees).